高瞻計畫_振動學課程 Lecture 1: Single Degree of Freedom Systems (III)

Prof. Kuo-Shen Chen
Department of Mechanical Engineering
National Cheng-Kung University

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Outline

- Rotordynamics, an introduction
- Impulse Responses
- Arbitrary excitation
- Transfer function and Laplace Domain
- Shock Isolation
- Simple problems
- Youtube Demos

Part I: Rotordynamics: an Introduction

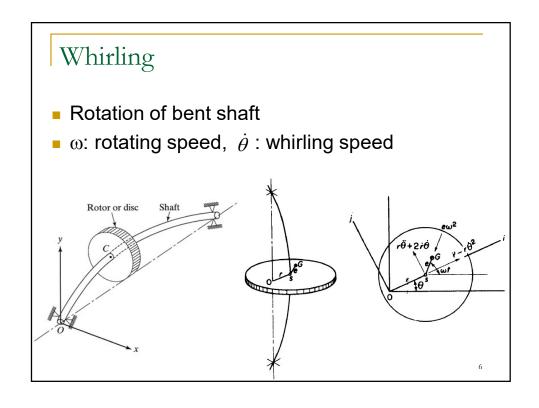
- Rotor phenomenon
- Unbalance
- Campbell diagrams

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Rotordynamics

- A specialized branch concerned with the behavior and diagnosis of rotating structures.
 - commonly used to analyze the behavior of structures ranging from <u>jet engines</u> and <u>steam turbines</u> to auto engines and computer <u>disk storage</u>.
 - Vibration, noise, bearing damages
- Key issues to be introduced
 - Critical speed
 - Whirling
 - Campbell diagram

Jeffcott Rotors A single disk mounted on a flexible and massless shaft with rigid bearings Serves as the fundamental model for studying rotor phenomena Flexible, Massless Shaft Disk Rigid Ks Massless Shaft Disk Figure 3.1-1 A simple Laval-Jeffcott rotor



Critical Speeds

• When frequency of rotation of shaft = one of the natural frequencies of the shaft, critical speed of undamped system:

$$\omega_n = \sqrt{\frac{k}{m}}$$

- When $\omega = \omega_n$, rotor undergoes large deflections
 - → cause fatique and damage bearings
- Slow transition of rotating shaft through the critical speed aids development of large amplitudes.
- Whirling critical speed should not be below 115 percent of the design full power speed (NOAA)

www.omao.noaa.gov/swath/contractdocs/attachments/AttachmentJ-02Rev2.pdf
Marine & Aviation Operaitons, National Oceanic & Atmospheric Administration, US Dept of Commerce

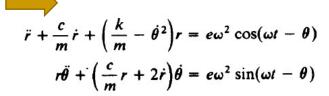
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Whirling: Mathematics

$$\mathbf{a}_{G} = \left[(\ddot{r} - r\dot{\theta}^{2}) - e\omega^{2} \cos(\omega t - \theta) \right] \mathbf{i} + \left[(r\ddot{\theta} + 2\dot{r}\dot{\theta}) - e\omega^{2} \sin(\omega t - \theta) \right] \mathbf{j}$$

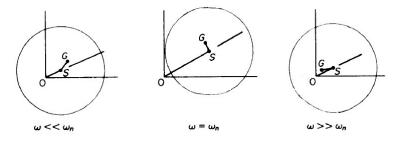
$$- kr - c\dot{r} = m \left[\ddot{r} - r\dot{\theta}^{2} - e\omega^{2} \cos(\omega t - \theta) \right]$$

$$- cr\dot{\theta} = m \left[r\ddot{\theta} + 2\dot{r}\dot{\theta} - e\omega^{2} \sin(\omega t - \theta) \right]$$



Whirling Behavior

• Synchronous whirling: $\dot{\theta} = \omega$

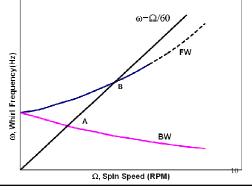


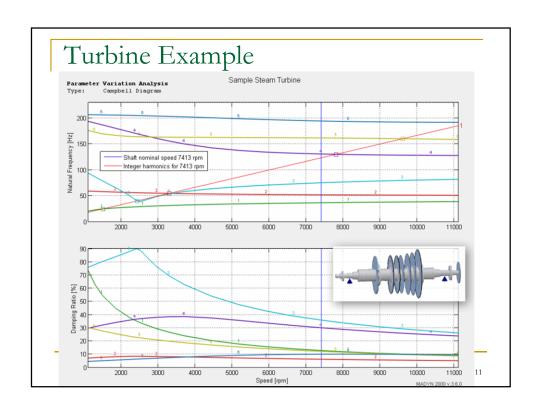
As rotating speed passes the critical speed, the imbalance location would actually move toward the center

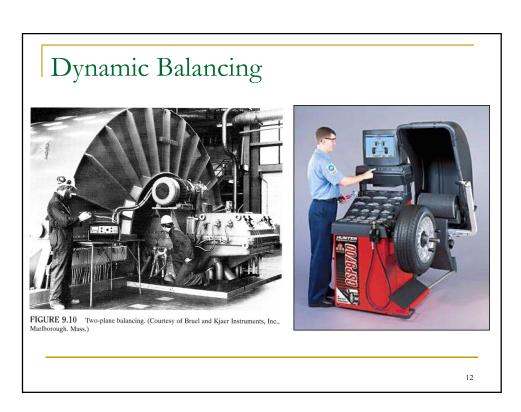
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Campbell Diagram

- known as "Whirl Speed Map" or a "Frequency Interference Diagram"
- Basic concept
 - Natural frequencies depends on rotating speed
 - Resonance occurs as natural frequencies hit rotating speed







Part II: Periodic Inputs

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Introduction

- In Lecture III, we have introduced the response of a SDOF system subjected to a single sinusoidal response
- How about the responses subjected to
 - □ A general periodic input
 - E.g., a saw tooth or a rectangular pulse train
 - A non-periodic input

Responses Under a General Force (I)

$$F(t) = \frac{a_0}{2} + \sum_{j=1}^{\infty} a_j \cos j\omega t + \sum_{j=1}^{\infty} b_j \sin j\omega t$$

$$a_j = \frac{2}{\tau} \int_0^{\tau} F(t) \cos j\omega t \, dt, \qquad j = 0, 1, 2, \dots$$

$$b_j = \frac{2}{\tau} \int_0^{\tau} F(t) \sin j\omega t \, dt, \qquad j = 1, 2, \dots$$

$$m\ddot{x} + c\dot{x} + kx = F(t) = \frac{a_0}{2} + \sum_{j=1}^{\infty} a_j \cos j\omega t + \sum_{j=1}^{\infty} b_j \sin j\omega t$$

$$m\ddot{x} + c\dot{x} + kx = \frac{a_0}{2}$$

$$m\ddot{x} + c\dot{x} + kx = a_j \cos j\omega t$$

$$m\ddot{x} + c\dot{x} + kx = b_j \sin j\omega t$$

Responses Under a General Force (II)

 $x_p(t) = \frac{a_0}{2L}$

$$\phi_{j} = \tan^{-1}\left(\frac{2\zeta jr}{1 - j^{2}r^{2}}\right)$$

$$r = \frac{\omega}{\omega_{n}}$$

$$x_{p}(t) = \frac{a_{0}}{2k} + \sum_{j=1}^{\infty} \frac{(a_{j}/k)}{\sqrt{(1 - j^{2}r^{2})^{2} + (2\zeta jr)^{2}}} \cos(j\omega t - \phi_{j})$$

$$+ \sum_{j=1}^{\infty} \frac{(b_{j}/k)}{\sqrt{(1 - j^{2}r^{2})^{2} + (2\zeta jr)^{2}}} \sin(j\omega t - \phi_{j})$$

Response Under a Periodic Force with Irregular Form

$$a_0 = \frac{2}{N} \sum_{i=1}^{N} F_i$$

$$a_j = \frac{2}{N} \sum_{i=1}^{N} F_i \cos \frac{2j\pi t_i}{\tau}, \qquad j = 1, 2, \dots$$

$$b_j = \frac{2}{N} \sum_{i=1}^{N} F_i \sin \frac{2j\pi t_i}{\tau}, \qquad j = 1, 2, \dots$$

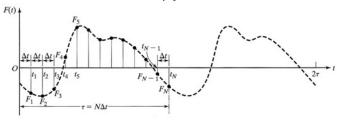


FIGURE 4.2 An irregular forcing function.

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Part III: Vibration Subject to Arbitrary Excitations

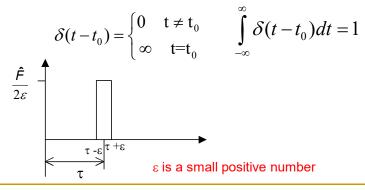
Introduction

- It is important to evaluate the response to a general, non-periodic input
- No, exact analytical solutions available
- However, the task can be performed by either
 - Convolution approach
 - □ Fourier transform or Laplace transform approach
 - Transfer function

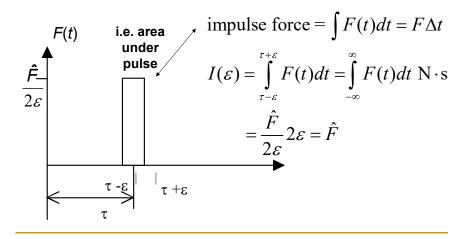
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Impulse Response

- The response of a vibration system subjected to a unit impulse input
- Impulse function



From sophomore dynamics:



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Equal

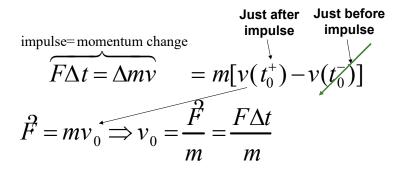
Use these properties to define the impulse function:

Dirac Delta

 $F(t) \qquad \text{function} \qquad \text{impulses}$ $F(t-\tau)=0, \quad t\neq \tau$ $\int_{-\infty}^{\infty} F(t-\tau)dt = \hat{F}$

If $\hat{F} = 1$, this is the Dirac Delta $\delta(t)$

Effect on spring-mass-damper?



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For an underdamped system

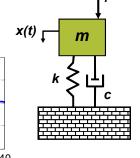
$$x(t) = \frac{\hat{F}e^{-\zeta\omega_n t}}{m\omega_d} \sin\omega_d t \text{ (response with zero I.C.)}$$

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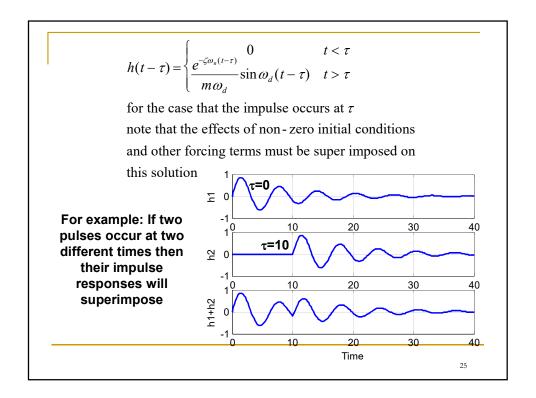
 $x(t) = \mathring{F}h(t)$, where $h(t) = \frac{e^{-\zeta \omega_n t}}{m\omega_d} \sin \omega_d t$

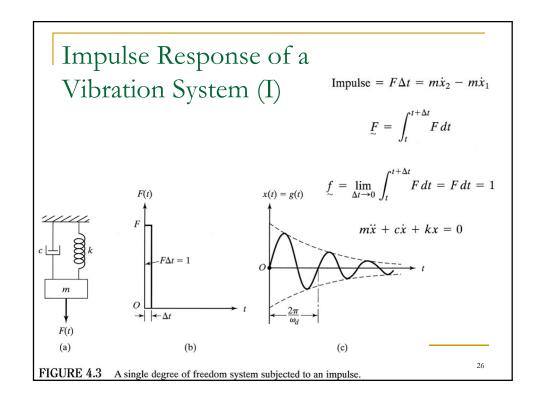
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0.5



 $\frac{h(t-\tau) = \frac{1}{m\omega_n} \sin \omega_n (t-\tau)}{m\omega_n}$





$$x(t) = e^{-\zeta \omega_n t} \left\{ x_0 \cos \omega_d t + \frac{\dot{x}_0 + \zeta \omega_n x_0}{\omega_d} \sin \omega_d t \right\}$$

$$\zeta = \frac{c}{2m\omega_n}$$

$$\omega_d = \omega_n \sqrt{1 - \zeta^2} = \sqrt{\frac{k}{m} - \left(\frac{c}{2m}\right)^2}$$

$$\omega_n = \sqrt{\frac{k}{m}}$$

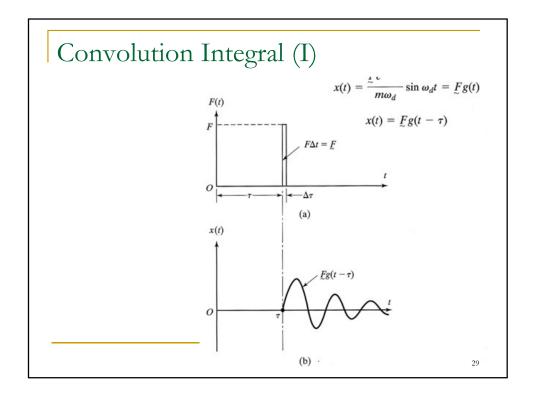
Impulse =
$$f = 1 = m\dot{x}(t = 0) - m\dot{x}(t = 0^{-}) = m\dot{x}_{0}$$

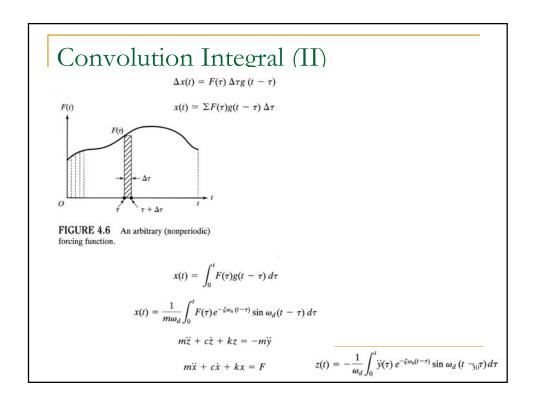
 $x(t = 0) = x_{0} = 0$
 $\dot{x}(t = 0) = \dot{x}_{0} = \frac{1}{m}$

$$x(t) = g(t) = \frac{e^{-\zeta \omega_n t}}{m\omega_d} \sin \omega_d t$$

Convolution Integrals: Introduction

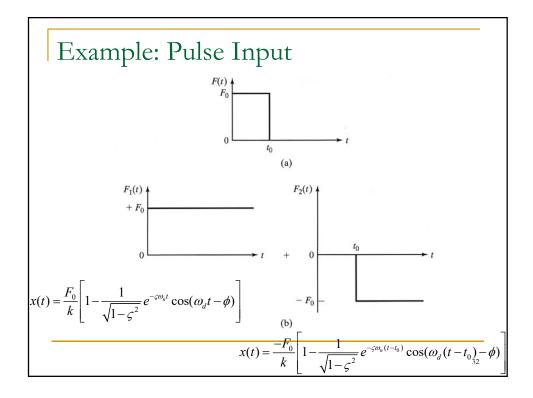
- Decompose a general input to the combination of a series of impulse train
- Superimpose the impulse response to form the final response
- The method itself can be treated as an "analytical" relation





Linear Superposition Method

- A complicate input can be decomposed into a few simple input with
 - scaling, multiplexing, and time shift operations
- The response are then becomes the linear superposition of the simple output after these linear operations



Time Shift Operation

- E.g., for the pulse input
 - \Box The total response = x1 + x2
 - Where
 - □ X1

$$x(t) = \frac{F_0}{k} \left[1 - \frac{1}{\sqrt{1 - \varsigma^2}} e^{-\varsigma \omega_n t} \cos(\omega_d t - \phi) \right]$$

□ x2

$$x(t) = \frac{-F_0}{k} \left[1 - \frac{1}{\sqrt{1 - \varsigma^2}} e^{-\varsigma \omega_n (t - t_0)} \cos(\omega_d (t - t_0) - \phi) \right]$$

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Part IV: Transfer Function and Laplace Transform

Transfer Functions

- A black box approach to correlate the inputoutput relation
- Based on linear superposition principle
- Usually performed by Laplace Transform approach
 - Or Fourier Transform

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Laplace Transform

- Changes ODE into algebraic equation
- Solve algebraic equation then compute the inverse transform
- Rule and table based in many cases
- Is used extensively in control analysis to examine the response
- Related to the frequency response function

$$X(s) = \mathcal{L}(x(t)) = \int_{0}^{\infty} x(t)e^{-st}dt$$

Take the transform of the equation of motion:

$$m\ddot{x} + c\dot{x} + kx = F_0 \cos \omega t \Rightarrow$$

 $(ms^2 + cs + k)X(s) = \frac{F_0 s}{s^2 + \omega^2}$

Now solve algebraic equation in s for X(s)

$$X(s) = \frac{F_0 s}{(ms^2 + cs + k)(s^2 + \omega^2)}$$

To get the time response this must be "inverse transformed"

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Laplace Transform: Fundamentals (I)

Consider a SDOF vibration system

$$m\ddot{x} + c\dot{x} + kx = F(t)$$

$$\overline{x}(s) = \mathcal{L}x(t) = \int_0^\infty e^{-st}x(t) dt$$

$$\mathcal{L}\frac{dx}{dt}(t) = e^{-st}x(t)\bigg|_{0}^{\infty} + s\int_{0}^{\infty} e^{-st}x(t) dt = s\overline{x}(s) - x(0)$$

$$\mathcal{L}\frac{d^2x}{dt^2}(t) = \int_0^\infty e^{-st} \frac{d^2x}{dt^2}(t) dt = s^2 \overline{x}(s) - sx(0) - \dot{x}(0)$$

$$\overline{F}(s) = \mathcal{L}F(t) = \int_0^\infty e^{-st} F(t) dt$$

Laplace Transform: Fundamentals (II)

$$(ms^{2} + cs + k) \overline{x}(s) = \overline{F}(s) + m\dot{x}(0) + (ms + c)x(0)$$

$$\overline{Z}(s) = \frac{\overline{F}(s)}{\overline{x}(s)} = ms^{2} + cs + k$$

$$\overline{Y}(s) = \frac{1}{\overline{Z}(s)} = \frac{\overline{x}(s)}{\overline{F}(s)} = \frac{1}{ms^{2} + cs + k} = \frac{1}{m(s^{2} + 2\zeta\omega_{n}s + \omega_{n}^{2})}$$

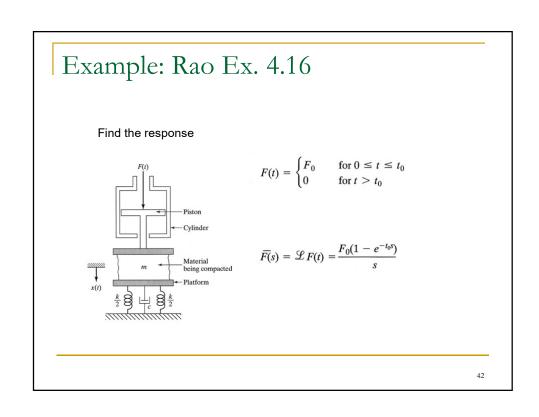
$$\overline{x}(s) = \overline{Y}(s)\overline{F}(s)$$

$$x(t) = \mathcal{L}^{-1}\overline{x}(s) = \mathcal{L}^{-1}\overline{Y}(s)\overline{F}(s)$$

Transfer Functions: Z(s) Impedance; Y(s): Admittance

	F(s)	$f(t), t \ge 0$	
Typical	1. 1	$\delta(t)$, unit impulse at $t=0$ $u_x(t)$, unit step	
	2. 1/s		
Laplace	3. n! sn+1	t"	
Transforms	$4. \frac{1}{s+s}$	e ^{-st}	
	$5. \frac{1}{(s+a)^n}$	$\frac{1}{(n-1)!}t^{n-1}e^{-at}$	
	6. $\frac{a}{s(s+a)}$	$1 - e^{-at}$	
	$7. \frac{1}{(s+a)(s+b)}$	$\frac{1}{b-a}(e^{-at}-e^{-bt})$	
	8. $\frac{s+\rho}{(s+a)(s+b)}$	$\frac{1}{b-a}[(p-a)e^{-at}-(p-b)e^{-bt}]$	
	$9. \frac{1}{(s+a)(s+b)(s+c)}$	$\frac{e^{-st}}{(b-s)(c-s)} + \frac{e^{-st}}{(c-b)(s-b)} + \frac{e^{-ct}}{(s-c)(b-c)}$	
	$10. \frac{s+p}{(s+a)(s+b)(s+c)}$	$\frac{(p-a)e^{-at}}{(b-a)(c-a)} + \frac{(p-b)e^{-bt}}{(c-b)(a-b)} + \frac{(p-c)e^{-ct}}{(a-c)(b-c)}$	
	$11. \frac{b}{s^2 + b^2}$	sin bt	
	12. $\frac{s}{s^2 + b^2}$	cos bt	
	13. $\frac{b}{(s+a)^2+b^2}$	e ^{-at} sin bt	
	14. $\frac{1}{(s+a)^2+b^2}$	e ^{-st} cos bt	
	$\frac{15.}{s^2+2\zeta\omega_n s+\omega_n^2}$	$\frac{\omega_n}{\sqrt{1-\zeta^2}}e^{-\zeta\omega_n t}\sin\omega_n\sqrt{1-\zeta^2t} \qquad \zeta<1$	
	$16. \frac{\omega_n^2}{s(s^2 + 2\zeta \omega_n s + \omega_n^2)}$	$1 + \frac{\omega_n}{\sqrt{1 - \zeta^2}} e^{-\beta \omega_n t} \sin\left(\omega_n \sqrt{1 - \zeta^2} t + \phi\right) \qquad \zeta < 1$	
		$\phi = \tan^{-1} \frac{\sqrt{1 - \xi^2}}{\xi} + \pi \text{ (third quadrant)}^{40}$	

	f(t)	$F(s) = \int_0^\infty f(t)e^{-st} dt$
Properties	1. $af_1(t) + bf_2(t)$	$aF_1(s) + bF_2(s)$
of Laplace	$2.\frac{df}{dt}$	sF(s)-f(0)
1 *	$3. \frac{d^2f}{dt^2}$	$s^2F(s) - sf(0) - \frac{df}{dt}\Big _{t=0}$
Transform	4. $\frac{d^n f}{dt^n}$	$s^n F(s) - \sum_{k=1}^n s^{n-k} g_{k-1}$
		$g_{\kappa-1} = \frac{d^{\kappa-1}f}{dt^{\kappa-1}}\bigg _{t=0}$
	$5. \int_0^t f(t) dt$	$\frac{F(s)}{s} + \frac{h(0)}{s}$
		$h(0) = \int f(t) dt \big _{t=0}$
	$6. g(t) = \begin{cases} 0 & t < D \\ f(t-D) & t \ge D \end{cases}$	$G(s) = e^{-sD}F(s)$
	7. $e^{-at}f(t)$	F(s+a)
	8. <i>tf</i> (<i>t</i>)	$-\frac{dF(s)}{ds}$
	9. $f(t) = \int_{0}^{t} x(t-\tau)y(\tau) d\tau = \int_{0}^{t} y(t-\tau)x(\tau) d\tau$	F(s) = X(s)Y(s)
	$10. f(\infty) = \lim_{s \to 0} sF(s)$	
	$11. f(0+) = \lim_{s \to \infty} sF(s)$	41



$$\overline{x}(s) = \frac{\overline{F}(s)}{m(s^2 + 2\zeta\omega_n s + \omega_n^2)} + \frac{s + 2\zeta\omega_n}{s^2 + 2\zeta\omega_n s + \omega_n^2} x_0$$

$$+ \frac{1}{s^2 + 2\zeta\omega_n s + \omega_n^2} \dot{x}_0$$

$$\overline{x}(s) = \frac{F_0(1 - e^{-t_0 s})}{ms(s^2 + 2\zeta\omega_n s + \omega_n^2)} + \frac{s + 2\zeta\omega_n}{s^2 + 2\zeta\omega_n s + \omega_n^2} x_0$$

$$+ \frac{1}{s^2 + 2\zeta\omega_n + \omega_n^2} \dot{x}_0$$

$$= \frac{F_0}{m\omega_n^2} \frac{1}{s\left(\frac{s^2}{\omega_n^2} + \frac{2\zeta s}{\omega_n} + 1\right)} - \frac{F_0}{m\omega_n^2} \frac{e^{-t_0 s}}{s\left(\frac{s^2}{\omega_n^2} + \frac{2\zeta s}{\omega_n} + 1\right)}$$

$$+ \frac{x_0}{\omega_n^2} \frac{s}{\left(\frac{s^2}{\omega_n^2} + \frac{2\zeta s}{\omega_n} + 1\right)} + \left(\frac{2\zeta x_0}{\omega_n} + \frac{\dot{x}_0}{\omega_n^2}\right) \frac{1}{\left(\frac{s^2}{\omega_n^2} + \frac{2\zeta s}{\omega_n} + 1\right)} \frac{1}{43}$$

$$x(t) = \frac{F_0}{m\omega_n^2 \sqrt{1 - \zeta^2}} \left[-e^{-\zeta \omega_n t} \sin(\omega_n \sqrt{1 - \zeta^2} t + \phi_1) + e^{-\zeta \omega_n (t - t_0)} \sin \left\{ \omega_n \sqrt{1 - \zeta^2} (t - t_0) + \phi_1 \right\} \right]$$

$$- \frac{x_0}{\sqrt{1 - \zeta^2}} e^{-\zeta \omega_n t} \sin(\omega_n \sqrt{1 - \zeta^2} t - \phi_1)$$

$$+ \frac{(2\zeta \omega_n x_0 + \dot{x}_0)}{\omega_n \sqrt{1 - \zeta^2}} e^{-\zeta \omega_n t} \sin(\omega_n \sqrt{1 - \zeta^2} t)$$

 $\phi_1 = \cos^{-1}(\zeta)$

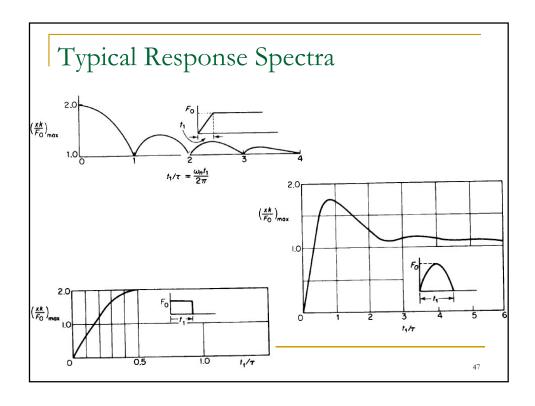
Part V: Response Spectrum

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Introduction

Shock

- A sudden application of a force input to a SDOF system to result a transient response
- The maximum value of the response can be used to measure the shook sensitivity
- Response spectrum is a plot of the maximum peak response of the SDOF oscillator as a function of natural frequency
- Different shock inputs result in different response spectra



General Response Spectra (1)

$$x(t) \bigg|_{\max} = \frac{1}{m\omega_n} \int_0^t F(\tau) \sin \omega_n (t - \tau) d\tau \bigg|_{\max}$$
$$S_d = \frac{S_v}{\omega_n}, \qquad S_a = \omega_n S_v$$

$$S_d = \frac{1}{\omega_n}, \qquad S_a = \omega_n S_v$$

$$\dot{z}(t) = -\frac{1}{\omega_d} \int_0^t \dot{y}(\tau) e^{-\zeta \omega_n (t-\tau)} \left[-\zeta \omega_n \sin \omega_d (t-\tau) + \omega_d \cos \omega_d (t-\tau) \right] d\tau$$

$$\dot{z}(t) = \frac{e^{-\zeta \omega_n t}}{\sqrt{1 - \zeta^2}} \sqrt{P^2 + Q^2} \sin(\omega_d t - \phi)$$

General Response Spectra (2)

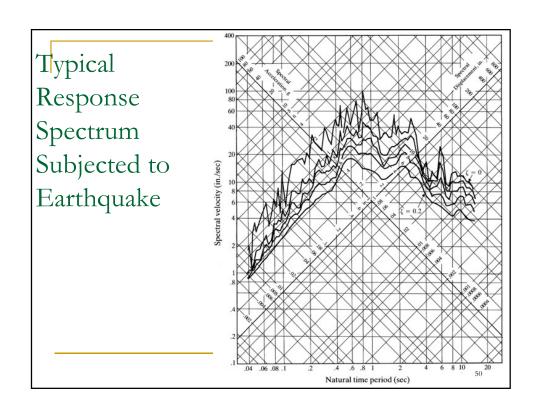
$$P = \int_0^t \ddot{y}(\tau) e^{\zeta \omega_n t} \cos \omega_d \tau d\tau$$

$$Q = \int_0^t \ddot{y}(\tau) e^{\zeta \omega_n t} \sin \omega_d \tau d\tau$$

$$\phi = \tan^{-1} \left\{ \frac{-(P\sqrt{1-\zeta^2} + Q\zeta)}{(P\zeta - Q\sqrt{1-\zeta^2})} \right\}$$

$$S_v = |\dot{z}(t)|_{\max} = \left| \frac{e^{-\zeta \omega_n t}}{\sqrt{1-\zeta^2}} \sqrt{P^2 + Q^2} \right|_{\max}$$

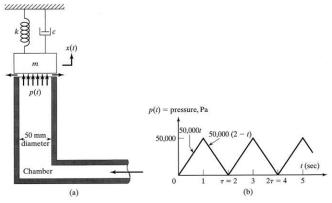
$$S_d = |z|_{\max} = \frac{S_v}{\omega_n}; \qquad S_v = |\dot{z}|_{\max}; \qquad S_a = |\ddot{z}|_{\max} = \omega_n S_v$$



Part VI: Simple Problems

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Problem 1. Periodic Vibration of a Hydraulic Valve (Rao 4.1)



In the study of vibrations of valves used in hydraulic control systems, the valve and its elastic stem are modeled as a damped spring-mass system, as shown in Fig. 4.1(a). In addition to the spring force and damping force, there is a fluid pressure force on the valve that changes with the amount of opening or closing of the valve. Find the steady-state response of the valve when the pressure in the chamber varies as indicated in Fig. 4.1(b). Assume k = 2500 N/m, c = 10 N-s/m, and m = 0.25 kg. 52

Problem 2. Response of a Structure under

Impact (Rao 4.4, 4.5)

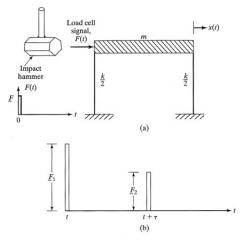
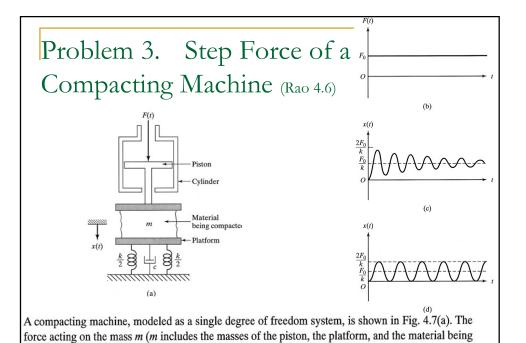


FIGURE 4.5 Structural testing using an impact hammer.

In the vibration testing of a structure, an impact hammer with a load cell to measure the impact force is used to cause excitation, as shown in Fig. 4.5(a). Assuming m = 5 kg, k = 2000 N/m c = 10 N-s/m and F = 20 N-s, find the response of the system.



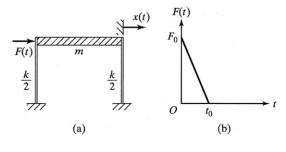
compacted) due to a sudden application of the pressure can be idealized as a step force, as shown in

Fig. 4.7(b). Determine the response of the system.

Problem 4. Blast Load on a Building

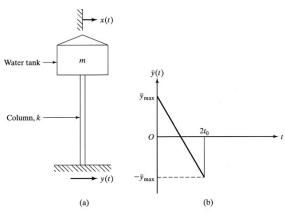
Frame (Rao 4.10)

$$F(\tau) = F_0 \left(1 - \frac{\tau}{t_0} \right) \text{ for } 0 \le \tau \le t_0$$
$$F(\tau) = 0 \quad \text{for } \tau > t_0$$



A building frame is modeled as an undamped single degree of freedom system (Fig. 4.11a). Find the response of the frame if it is subjected to a blast loading represented by the triangular pulse shown in Fig. 4.11(b).

Problem 5: Water Tank Subjected to Base Acceleration (Rao 4.12)



The water tank, shown in Fig. 4.13(a), is subjected to a linearly varying ground acceleration as shown in Fig. 4.13(b) due to an earthquake. The mass of the tank is m, the stiffness of the column is k, and damping is negligible. Find the response spectrum for the relative displacement, $z = x_0 - y$, of the water tank.

Problem 6. Response of a Building Frame to an Earthquake (Rao. 4.13)

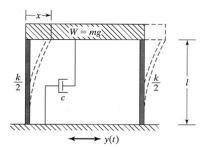


FIGURE 4.17 Building frame subjected to base motion.

A building frame has a mass of 6,800 kg and two columns of total stiffness k, as indicated in Fig. 4.17. It has a damping ratio of 0.05 and a natural time period of 1.0 sec. For the earthquake characterized in Fig. 4.15, determine the following:

Part VII: Youtube Demonstrations